

$$f_0 p_0 - f_1 \left(p_0 + \frac{\partial p}{\partial x} dx \right) - dR = dm \frac{dc}{d\tau}$$

with $dx \rightarrow 0$: $f_0 \rightarrow f_1 \rightarrow f$

$$-f \frac{\partial p}{\partial x} dx - dR = dm \frac{dc}{d\tau} \quad \Bigg| : dm = \frac{f dx}{v}$$

$$-v \frac{\partial p}{\partial x} - R = \frac{dc}{d\tau}$$

wh $R = \frac{dR}{dm}$ is resistance force per 1 kg of flowing gas.

er
For steady state mode

$$\frac{\partial p}{\partial x} = \frac{dp}{dx}$$

$$-v dp - R dx = \frac{dx}{d\tau} dc \quad \xrightarrow{\frac{dx}{d\tau} = c} \quad \boxed{cdc = -v dp - R dx}$$

The equation of the change in the amount of flow in unidirectional flow

Integrating from 0-0 to 1-1

$$\frac{c_1^2 - c_0^2}{2} = \int_{p_1}^{p_0} v dp - \int_{x_0}^{x_1} R dx$$

$$\frac{c_1^2 - c_0^2}{2} = \int_{p_1}^{p_0} \nu dp - \int_{x_0}^{x_1} R dx$$

* Theoretical expansion

$$R(x) = 0$$

$$\frac{c_{1t}^2 - c_0^2}{2} = \int_{p_1}^{p_0} \nu dp$$

$$\int_{p_1}^{p_0} \nu dp = l_{mexH} = h_0 - h_{1t}$$

$$\frac{c_{1t}^2 - c_0^2}{2} = h_0 - h_{1t}$$

** Real expansion

$$R = f(x)$$

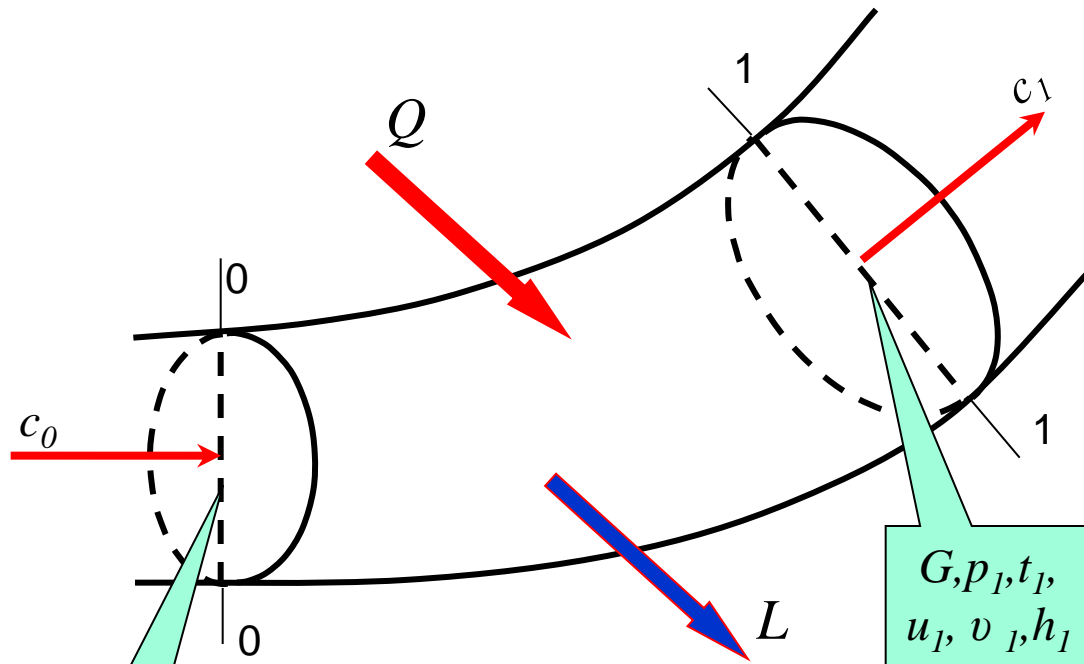
It is known that $R = f(\Delta, Re = \varphi(c, \nu, d), \text{ channel curvature, ...})$

in particular, the possibility of separation of the boundary layer under diffuse flow

In the turbine theory, the momentum equation for the flow in the channel in case of actual flow is replaced with the experimental data.

2.1.4. Energy-conservation equation

(for the flow)



$G, p_0, t_0,$
 u_0, v_0, h_0

$G, p_1, t_1,$
 u_1, v_1, h_1

Particular cases

a) $q = 0$

$$\frac{c_0^2}{2} + h_0 = \frac{c_1^2}{2} + h_1 + l$$

b) $q = 0, l = 0$

$$\frac{c_0^2}{2} + h_0 = \frac{c_1^2}{2} + h_1$$

Supplied energy

$$\frac{Gc_0^2}{2} + Gh_0 + Q$$

Withdrawn energy

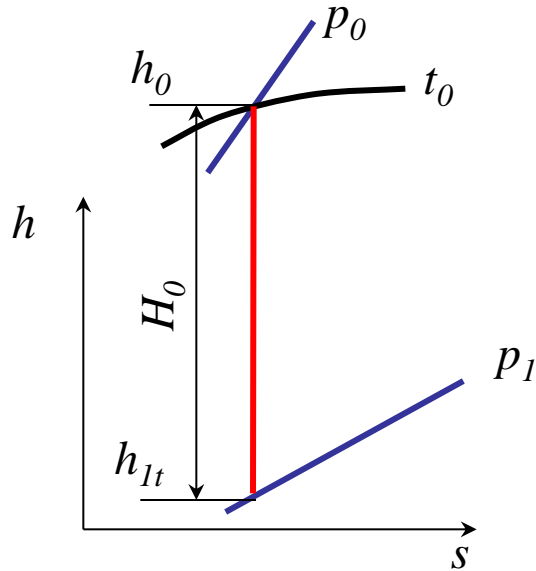
$$\frac{Gc_1^2}{2} + Gh_1 + L$$

$$\frac{c_0^2}{2} + h_0 + q = \frac{c_1^2}{2} + h_1 + l$$

p_0, t_0, p_1 are set. Determine the velocity at the channel outlet when

$$q = 0, l = 0$$

* Theoretical expansion



$$\frac{c_0^2}{2} + h_0 = \frac{c_{1t}^2}{2} + h_{1t} \quad \frac{c_{1t}^2 - c_0^2}{2} = h_0 - h_{1t}$$

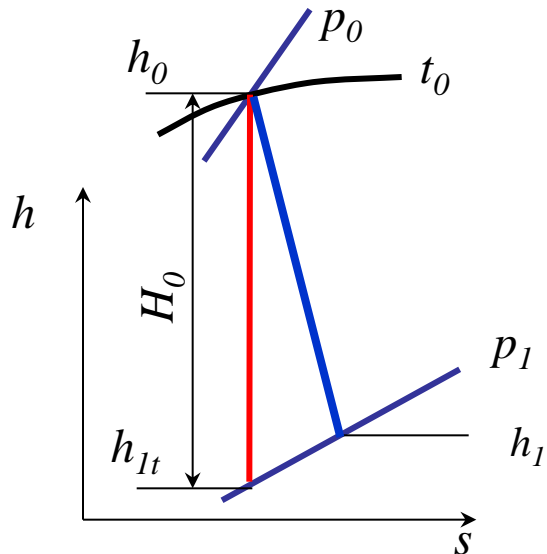
$$c_{1t} = \sqrt{2(h_0 - h_{1t}) + c_0^2} \quad \leftarrow \text{h в Дж/кг}$$

If h in kJ/kg $\implies c_{1t} = \sqrt{2000(h_0 - h_{1t}) + c_0^2}$

$$c_{1t} = 44,7 \sqrt{(h_0 - h_{1t}) + \frac{c_0^2}{2000}}$$

$H_0 = h_0 - h_{1t}$ is available heat drop per channel

** Real expansion



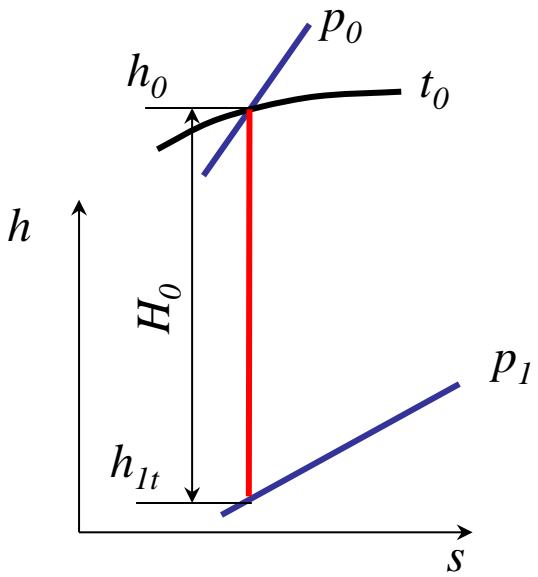
$$\frac{c_0^2}{2} + h_0 = \frac{c_1^2}{2} + h_1$$

$$c_1 = 44,7 \sqrt{(h_0 - h_1) + \frac{c_0^2}{2000}}$$

$H_i = h_0 - h_1$ is real heat drop per channel

Comparison of application of the equation of momentum and energy conservation equation

Theoretical expansion



The equation of the amount of motion

$$\frac{c_{1t}^2 - c_0^2}{2} = h_0 - h_{1t}$$

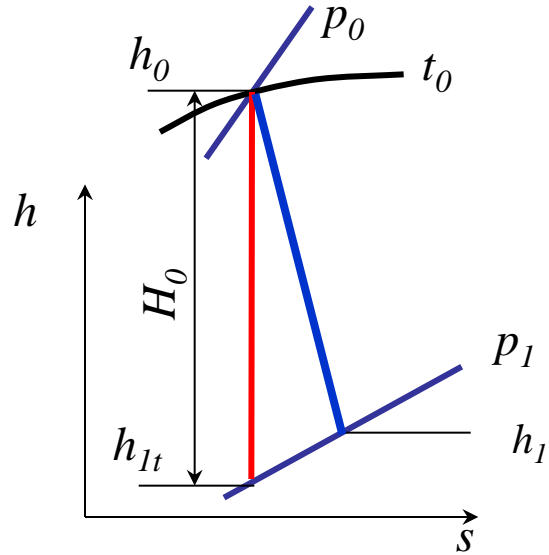
$$c_{1t} = \sqrt{2(h_0 - h_{1t}) + c_0^2}$$

Energy conservation equation

$$\frac{c_{1t}^2 - c_0^2}{2} = h_0 - h_{1t}$$

$$c_{1t} = \sqrt{2(h_0 - h_{1t}) + c_0^2}$$

Real expansion



$$\frac{c_1^2 - c_0^2}{2} = \int_{p_1}^{p_0} v dp - \int_{x_0}^{x_1} R dx$$

It is necessary to know the thermodynamic process along the channel

$$\frac{c_1^2 - c_0^2}{2} = h_0 - h_1$$

$$c_1 = \sqrt{2(h_0 - h_1) + c_0^2}$$

It is necessary to know the initial and final state

*** Flow acceleration

$$\frac{c_1^2 - c_0^2}{2} = h_0 - h_1$$

If $h_1 < h_0$, $c_1 > c_0$ the flow **accelerates** (convergent flow)

If $h_1 > h_0$, $c_1 < c_0$ the flow **decelerates** (divergent flow)

2.2. The flow characteristics during isentropic expansion of gas in the channels

The channel, in which the flow smoothly accelerates, is called a **nozzle one** or a **nozzle**.

The channel, in which the flow smoothly slows down, is called a **diffusion one** or a **diffuser**.